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Section 1

In this note we will compute the correlation between the beam parameters at the output of a magnetic channel and the errors due to the channel assembling and to the variations of current supply.

In other words we will compute the effects of a displacement of any element on the final beam features.

Mathematically speaking this computation is equivalent to a computation of derivatives of certain quantities with respect to some parameters.

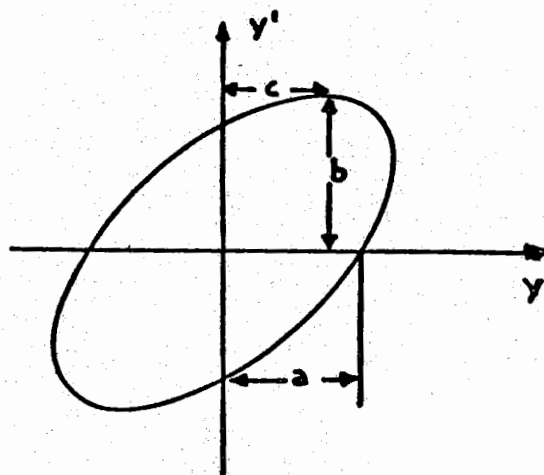
For each of the two phase planes we will consider the following final beam parameters:

- a) - center coordinates of the ellipse that envelopes the beam, Y , Y' ;
- b) - R , X defined as: $R = a/b$; $X = c/b$ (see fig. 1);
- c) - absolute maximum values of the displacement and the inclination, for a non centered ellipse, defined as:

$$(1) \quad \begin{aligned} Y_{\max} &= \sqrt{E(R+X^2/R)} + |Y| \\ Y'_{\max} &= \sqrt{E/R} + |Y'| \end{aligned}$$

2.

where $E = ab$ is the emittance.



$$R = \frac{a}{b}$$

$$X = \frac{c}{b}$$

FIG. 1

Likewise the parameters variations we consider are:

- d) - variations of the drift spaces lengths;
- e) - variations of the quadrupole lenses or bending magnets current supplies;
- f) - rigid displacements of the quadrupoles or magnets, that we will subdivide into translations and rotations.

To compute the effects of the above mentioned errors we will proceed in the following way:

making use of the well-known matrix method⁽¹⁾ the channel is first traced assuming magnetic elements without errors so that we can obtain the theoretical values of the variables defined in a), b), c).

Successively the channel is traced many times, every time introducing a small but finite variation (such as in d), e), f)) for the single element.

Such variations are obtained by suitably modifying the computation corresponding to the element we are considering.

In such a way we obtain the varied final parameters, which together with the theoretical ones allow us to compute the derivatives of the final parameters with respect to the varied parameters (properly speaking such derivatives are "incremental ratios").

The method we will use to treat the rigid displacements is too rigorous and not consistent with the first order treatment of the channel, but we have adopted it for two reasons:

- a) - It is easier to draw a rigorous formula than a good approximation;

- b) - The formulae we have deduced are true also if the analytical treatment of the channel is computed in a higher approximation.

Section 2

The computation of the effects caused by the variations we mentioned in e), d) is easily performed because it involves only the change of the magnetic element matrix.

Thus the variation of a drift space length only produces in the corresponding matrix the simple change:

$$\begin{vmatrix} 1 & \ell \\ 0 & 1 \end{vmatrix} \longrightarrow \begin{vmatrix} 1 & \ell + \Delta\ell \\ 0 & 1 \end{vmatrix}$$

A variation of the current supply of a magnetic element is equivalent for relativistic particles, to a percentual variation of the particle energy (because of the Lorentz force) and in this case the matrices are easily modified.

Properly speaking if the current supply i goes into $i(1 + \Delta i/i)$ the constant k of the quadrupole matrix becomes $k(1 + \Delta i/i)$ because it is proportional to the magnetic field; whereas in the third-order bending matrices, one must introduce a $\Delta p/p = \Delta i/i$.

Section 3

On the contrary the computation of the effects caused by rigid displacements of magnetic elements, is more difficult.

To define such displacements we shall assume a right-handed reference system solidary with the element. We shall also assume that the origin of the systems is coincident, for quadrupole, with its central point. For a bending magnet the origin of the reference is taken at the central point of the principal orbit. The reference axes are oriented as in fig. 2.

As shown in fig. 2 we call T_1 the reference placed at the input face of the element; T_2 the central reference, and T_3 the one placed at the output face of the element.

For what follows it is useful to compute the coordinates of the origin and of the unit vectors of T_1 (T_3) with respect to T_2 .

4.

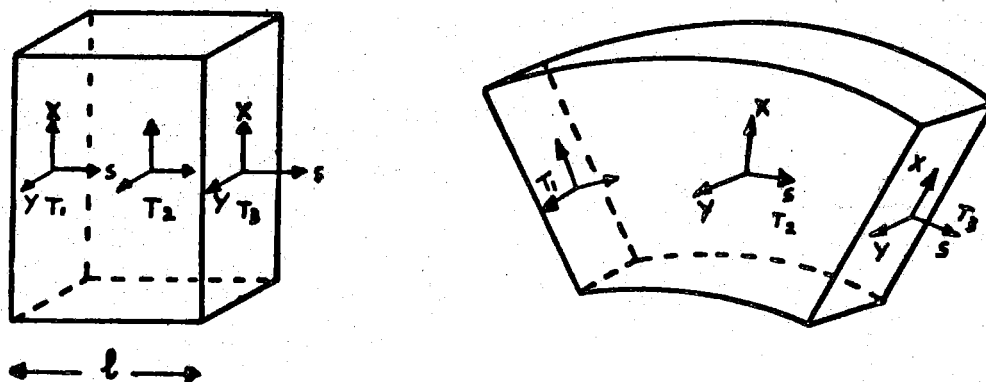


FIG. 2

We thus obtain for each kind of element (see fig. 3):

1) - Quadrupole - (We indicate the element length with l)

$$x_A = 0.$$

$$x_B = 1.$$

$$y_A = 0.$$

$$y_B = 0.$$

$$s_A = -.5 \cdot l$$

$$s_B = -.5 \cdot l$$

(2)

$$x_C = 0.$$

$$x_D = 0.$$

$$y_C = 1.$$

$$y_D = 0.$$

$$s_C = -.5 \cdot l$$

$$s_D = -.5 \cdot l + 1$$

2) - Magnet - (We indicate the deflection angle with α and the radius of curvature with r)

$$x_A = -r(1 - \cos \frac{\alpha}{2})$$

$$x_B = -r(1 - \cos \frac{\alpha}{2}) + \cos \frac{\alpha}{2}$$

$$y_A = 0.$$

$$y_B = 0.$$

$$s_A = -r \sin \frac{\alpha}{2}$$

$$s_B = -r \sin \frac{\alpha}{2} - \sin \frac{\alpha}{2}$$

(2')

$$x_C = -r(1 - \cos \frac{\alpha}{2})$$

$$x_D = -r(1 - \cos \frac{\alpha}{2}) + \sin \frac{\alpha}{2}$$

$$y_C = 1.$$

$$y_D = 0.$$

$$s_C = -r \sin \frac{\alpha}{2}$$

$$s_D = -r \sin \frac{\alpha}{2} + \cos \frac{\alpha}{2}$$

To obtain the coordinates of the same points with respect to system T_3 it is sufficient to change, for the quadrupoles, the sign of ν , and, for the bending magnets the sign of α .

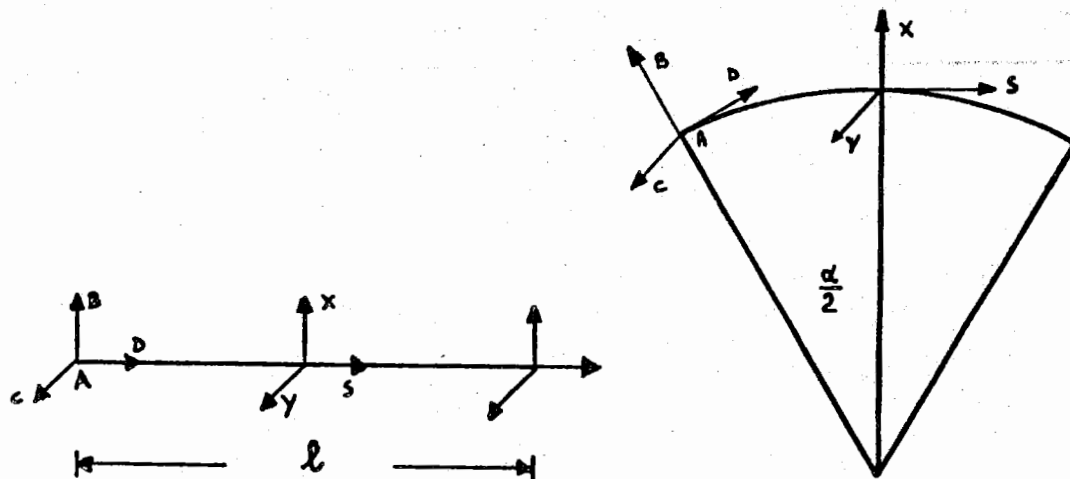


FIG. 3

Section 4

With \bar{T}_1 , \bar{T}_2 , \bar{T}_3 we indicate the coordinate systems of the displaced elements.

We define the displacement of the generic element giving the \mathcal{Z} transformation of the displacement of \bar{T}_2 with respect to T_2 , i. e. giving three translation parameters p_* , q_* , t_* , along the three axes, and nine direction cosines forming the rotation matrix:

$$R = \begin{vmatrix} \alpha_1 & \alpha_2 & \alpha_3 \\ \delta_1 & \delta_2 & \delta_3 \\ \psi_1 & \psi_2 & \psi_3 \end{vmatrix}$$

Such parameters have the following meaning:

p_*	q_*	t_*	center coordinates of \bar{T}_2 with respect to T_2
α_1	α_2	α_3	projections of the three unitary vectors of \bar{T}_2 onto the x axis of T_2
δ_1	δ_2	δ_3	projections of the three unitary vectors of \bar{T}_2 onto the y axis of T_2

6.

$\varphi_1 \quad \varphi_2 \quad \varphi_3$ projections of the three unitary vectors of \overline{T}_2 onto the z axis of T_2

We will consider the following kinds of displacements of \overline{T}_2 with respect to T_2 :

A) - Rotations

In this case one always has

$$p_{\star} = 0 \quad q_{\star} = 0 \quad t_{\star} = 0$$

and also:

1) rotation through an angle φ_x about the x axis:

$$R = \begin{vmatrix} 1 & 0 & 0 \\ 0 & \cos \varphi_x & -\text{sen } \varphi_x \\ 0 & \text{sen } \varphi_y & \cos \varphi_y \end{vmatrix}$$

2) rotation through an angle φ_y about the y axis:

$$R = \begin{vmatrix} \cos \varphi_y & 0 & \text{sen } \varphi_y \\ 0 & 1 & 0 \\ -\text{sen } \varphi_y & 0 & \cos \varphi_y \end{vmatrix}$$

3) rotation through an angle φ_s about the s axis:

$$R = \begin{vmatrix} \cos \varphi_s & -\text{sen } \varphi_s & 0 \\ \text{sen } \varphi_s & \cos \varphi_s & 0 \\ 0 & 0 & 1 \end{vmatrix}$$

B) - Translation

one always has

$$R = \begin{vmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{vmatrix}$$

and also:

4) translation through a length Δx along x axis

$$p_{\star} = \Delta x \quad q_{\star} = 0 \quad t_{\star} = 0$$

5) translation through a length Δy along y axis

$$p_{\star} = 0 \quad q_{\star} = \Delta y \quad t_{\star} = 0$$

6) translation through a length Δs along s axis

$$p_{\star} = 0 \quad q_{\star} = 0 \quad t_{\star} = \Delta s$$

Section 5

Since we are dealing with rigid displacements, when the parameters defining the displacement of T_2 have been assigned the ones \bar{T}_1 ; \bar{T}_3 are also defined.

In the present section we will calculate the transformation law between the initial systems T_1 and \bar{T}_1 .

From the above considerations it is clear that when the element considered is displaced, the four points A, B, C, D go into the four points A', B', C', D' whose coordinates with respect to T_2 can be calculated by the \mathcal{Z} transformation defined in section 4.

If, for example, we apply \mathcal{Z} to A we have:

$$(3) \quad \begin{aligned} x_{A'} &= p_{\star} + \mathcal{V}_1 x_A + \mathcal{V}_2 y_A + \mathcal{V}_3 s_A \\ y_{A'} &= q_{\star} + \mathcal{D}_1 x_A + \mathcal{D}_2 y_A + \mathcal{D}_3 s_A \\ s_{A'} &= t_{\star} + \mathcal{Y}_1 x_A + \mathcal{Y}_2 y_A + \mathcal{Y}_3 s_A \end{aligned}$$

Likewise for B', C', D'.

Let us consider now the points A, B, C, D, A', B', C', D', as vectors whose components are $x_A, y_A, s_A; x_B, y_B, s_B; \dots; x_{A'}, y_{A'}, s_{A'}; \dots$ and let us write them as $\vec{A}, \vec{B}, \vec{C}, \vec{A}', \vec{B}', \vec{C}', \dots$

By means of suitable differences we obtain the following vectors:

$$(4) \quad \begin{aligned} (\vec{B} - \vec{A}) & \text{ unitary vector of } T_1 \text{ x axis} \\ (\vec{C} - \vec{A}) & \text{ " " " " y "} \\ (\vec{D} - \vec{A}) & \text{ " " " " s "} \\ (\vec{A}' - \vec{A}) & \text{ difference between the } \bar{T}_1 \text{ and } T_1 \text{ origins} \\ (\vec{B}' - \vec{A}') & \text{ unitary vector of x axis of } \bar{T}_1 \\ (\vec{C}' - \vec{A}') & \text{ " " " y " " "} \\ (\vec{D}' - \vec{A}') & \text{ " " " s " " "} \end{aligned}$$

8.

Let us now consider as un-known terms the parameters of the transformation \mathcal{F} that makes us pass from the coordinates $\bar{x}, \bar{y}, \bar{s}$ of \bar{T}_1 to the coordinates x, y, s of T_1 of a point P.

$$(5) \quad \begin{aligned} x &= x_{\star} + \alpha_1 \bar{x} + \beta_1 \bar{y} + \gamma_1 \bar{s} \\ y &= y_{\star} + \alpha_2 \bar{x} + \beta_2 \bar{y} + \gamma_2 \bar{s} \\ s &= s_{\star} + \alpha_3 \bar{x} + \beta_3 \bar{y} + \gamma_3 \bar{s} \end{aligned}$$

keeping in mind the meaning of (5) and the above given definitions (4), we obtain:

$$\begin{aligned} x_{\star} &= (\vec{A}' - \vec{A}) \cdot (\vec{B} - \vec{A}) & \alpha_1 &= (\vec{B}' - \vec{A}') \cdot (\vec{B} - \vec{A}) \\ y_{\star} &= (\vec{A}' - \vec{A}) \cdot (\vec{C} - \vec{A}) & \alpha_2 &= (\vec{B}' - \vec{A}') \cdot (\vec{C} - \vec{A}) \\ s_{\star} &= (\vec{A}' - \vec{A}) \cdot (\vec{D} - \vec{A}) & \alpha_3 &= (\vec{B}' - \vec{A}') \cdot (\vec{D} - \vec{A}) \\ \beta_1 &= (\vec{C}' - \vec{A}') \cdot (\vec{B} - \vec{A}) & \gamma_1 &= (\vec{D}' - \vec{A}') \cdot (\vec{B} - \vec{A}) \\ \beta_2 &= (\vec{C}' - \vec{A}') \cdot (\vec{C} - \vec{A}) & \gamma_2 &= (\vec{D}' - \vec{A}') \cdot (\vec{C} - \vec{A}) \\ \beta_3 &= (\vec{C}' - \vec{A}') \cdot (\vec{D} - \vec{A}) & \gamma_3 &= (\vec{D}' - \vec{A}') \cdot (\vec{D} - \vec{A}) \end{aligned}$$

The indicated inner products can easily be computed in the T_2 system using the coordinates as in (2) (2') (3').

Let us remark that the \mathcal{C} transformation connects the coordinates of two different points A, A' referred to the same frame T_2 ; the \mathcal{F} transformation connects the coordinates of a same point P referred to two different frames (T_1, \bar{T}_1).

Section 6

In this section our task is to calculate the parameters of the generic particle at the input face of the displaced element in reference T_1 .

The particle before entering a magnetic element follows a straight path (assuming the rectangular model, as we do) and thus we can write the trajectory equation in T_1 as:

$$(6) \quad \begin{aligned} x &= x_0 + x'_0 s \\ y &= y_0 + y'_0 s \end{aligned}$$

where (x_0, x'_0, y_0, y'_0) are the four trajectory parameters and s is the curvilinear abscissa.

If the center of reference \bar{T}_1 is displaced with respect to T_1 by a vector whose components are x_*, y_*, s_* and if $\sigma_1, \sigma_2, \sigma_3$ are the direction cosines of the ξ, η plane in T_1 the equation of this plane in T_1 is

$$(7) \quad \sigma_1(x - x_*) + \sigma_2(y - y_*) + \sigma_3(s - s_*) = 0$$

the intersection of the straight line (6) with the plane (7) gives us the coordinates (referred to T_1) of the particle input in the displaced element.

Nevertheless in order to calculate the slopes of the tangents to the path in the displaced element it is more useful for us to write the intersection of (6) with a plane parallel to the ξ, η one, at a distance ζ , and that has equation:

$$(7') \quad \sigma_1(x - x_*) + \sigma_2(y - y_*) + \sigma_3(s - s_*) = \zeta$$

we thus obtain:

$$(8) \quad x = \frac{x_0(\sigma_3 + \sigma_2 y'_0) - y_0 \sigma_2 x'_0 + x'_0(\sigma_1 x_* + \sigma_2 y_* + \sigma_3 s_* + \zeta)}{\sigma_1 x'_0 + \sigma_2 y'_0 + \sigma_3}$$

$$y = \frac{y_0(\sigma_3 + \sigma_1 x'_0) - x_0 \sigma_1 y'_0 + y'_0(\sigma_1 x_* + \sigma_2 y_* + \sigma_3 s_* + \zeta)}{\sigma_1 x'_0 + \sigma_2 y'_0 + \sigma_3}$$

$$s = \frac{\sigma_1 x_* + \sigma_2 y_* + \sigma_3 s_* + \zeta - \sigma_1 x_0 - \sigma_2 y_0}{\sigma_1 x'_0 + \sigma_2 y'_0 + \sigma_3}$$

Let us now specify exactly the displacement we are considering.

Keeping in mind that the direction cosines of the unitary vectors of \bar{T}_1 have been labelled $\alpha_1, \alpha_2, \alpha_3, \beta_1, \beta_2, \beta_3, \bar{\sigma}_1, \bar{\sigma}_2, \bar{\sigma}_3$, in order to obtain the input coordinates ξ, η respect to T_1 we must perform the inner product of vector $\vec{i}(x-x_*) + \vec{j}(y-y_*) + \vec{k}(s-s_*)$ (i. e. vectorial distance between the computed intersection and the origin of T_1 with the unitary vector of the ξ axis of T_1 ($\vec{i} \alpha_1 + \vec{j} \alpha_2 + \vec{k} \alpha_3$),

10.

and with the unitary vector of the η axis of \bar{T}_1 ($\vec{i} \beta_1 + \vec{j} \beta_2 + \vec{k} \beta_3$) respectively.

We thus have:

$$\xi = \alpha_1(x - x_{\star}) + \alpha_2(y - y_{\star}) + \alpha_3(s - s_{\star})$$

$$\eta = \beta_1(x - x_{\star}) + \beta_2(y - y_{\star}) + \beta_3(s - s_{\star})$$

In order to compute the derivatives ξ' η' we must derive the eq. (9) with respect to ξ .

So doing we have:

$$(10) \quad \xi' = \frac{\alpha_1 x'_0 + \alpha_2 y'_0 + \alpha_3}{\sigma_1 x'_0 + \sigma_2 y'_0 + \sigma_3}$$

$$\eta' = \frac{\beta_1 x'_0 + \beta_2 y'_0 + \beta_3}{\sigma_1 x'_0 + \sigma_2 y'_0 + \sigma_3}$$

The formula (10) gives directly the slopes in the input face of the displaced element whereas to obtain ξ η we must introduce in formula (9) the x , y , s values given by (8) assuming a vanishing ξ .

Section 7

It is now interesting to compute the jacobian of the transformation corresponding to the input face of the displaced magnetic element.

$$\xi = \xi(x_0, y_0, x'_0, y'_0)$$

$$\eta = \eta(x_0, y_0, x'_0, y'_0)$$

$$\xi' = \xi'(x_0, y_0, x'_0, y'_0)$$

$$\eta' = \eta'(x_0, y_0, x'_0, y'_0)$$

We can thus compute by how much the jacobian differs from unity. Let us remark that the transformations (8), (9), (10) are non-linear.

Consider the usual expression:

$$J = \begin{vmatrix} \xi_{x_0} & \xi_{y_0} & \xi_{x'_0} & \xi_{y'_0} \\ \eta_{x_0} & \eta_{y_0} & \eta_{x'_0} & \eta_{y'_0} \\ \xi'_{x_0} & \xi'_{y_0} & \xi'_{x'_0} & \xi'_{y'_0} \\ \eta'_{x_0} & \eta'_{y_0} & \eta'_{x'_0} & \eta'_{y'_0} \end{vmatrix}$$

Where with the subscript we indicate the derivative. From (10) we deduce that the derivatives ξ' η' with respect to x_0 , y_0 are zero, and we can write:

$$J = (\xi_{x_0} \eta_{y_0} - \eta_{x_0} \xi_{y_0}) (\xi'_{x'_0} \eta'_{y'_0} - \eta'_{x'_0} \xi'_{y'_0})$$

After a rather cumbersome computation we have:

$$(\xi_{x_0} \eta_{y_0} - \eta_{x_0} \xi_{y_0}) = \frac{1}{(\sigma_3 + \sigma_2 y'_0 + \sigma_1 x'_0)}$$

$$(\xi'_{x'_0} \eta'_{y'_0} - \eta'_{x'_0} \xi'_{y'_0}) = \frac{1}{(\sigma_3 + \sigma_2 y'_0 + \sigma_1 x'_0)^3}$$

follows:

$$(11) \quad J = \frac{1}{(\sigma_3 + \sigma_2 y'_0 + \sigma_1 x'_0)^3}$$

Analogous formulae hold for the output face of the element.

Let us note that if in (11) we consider the amplitudes of the rotation angles as first order infinitesimals (like y'_0 and x'_0 in first order approximation) it follows that J , neglecting the second order terms, is equal to one consistently with the linear treatment of the channel.

Would actually be contradictory to take into consideration the second order infinitesimals at the input and at the output of the displaced elements and then to assume the determinants of the transfer matrices of such elements to be equal to one.

In any case the effect of (11) is negligible for a magnetic chan

nel if the particles go through the element only once.

Let us now remark that such a development is not applicable to the circular machines if one considers the second order terms.

It remains to be noted that the result (11) gives, is not generally speaking in contradiction with Liouville's theorem which deals with the whole phase space.

The problem becomes more interesting if we are dealing with a circular machine and if the effect caused by the errors is due to an arbitrary configuration of the magnetic field.

For example this case arises in the FFAG machines where the magnetic field, as compared to that of a traditional machine, is folded along spirals.

It is clear that in such case one is in need of a different mathematical procedure.

Two recent works^(2, 3) should be useful in this respect, the first concerning Liouville's theorem and the general problem of the stability, the second the motion of a charged particle in a magnetic field.

Section 8

In the last section we have computed the values of the particle-parameters at the input face of the displaced element.

Applying now the well known matrix-method we obtain the particle parameters at the output face of the element.

To apply the matrix of the subsequent drift space to the displaced magnetic element we must compute the particle path parameters in a P point of the x, y face of T_3 (see fig. 4).

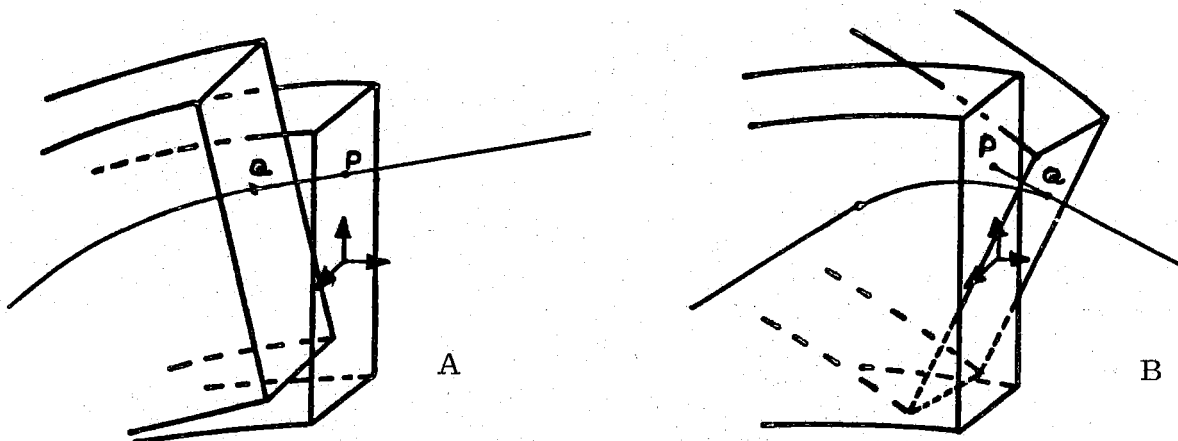


FIG. 4

As shown in fig. 4A and 4B it is clear that, in general, the point P of the x, y face of T_3 is not a point of the path; nevertheless P is always the intersection of the tangent to the trajectory at point Q of the displaced element, with the xy face of T_3 .

The same is true for the input face of the element, although the latter case looks easier. But such a difference is only apparent since reversing the direction of the particle on the trajectory the two cases are interchanged.

According to what we said the computation is very similar to the one of section 2 - 3.

More exactly to obtain the transformation that connects \overline{T}_3 to T_3 one must do all the computations of section 2 (as seen earlier it is enough to change the sign of ℓ for the quadrupoles and of α for the bending magnets).

Then one must compute the inverse transformation, i. e. the transformation that connects the coordinates $\overline{x}, \overline{y}, \overline{s}$ of \overline{T}_3 to the x, y, s of T_3 .

One obtains:

$$(5') \quad \begin{aligned} x &= p_{\star} + \alpha_1 x + \alpha_2 y + \alpha_3 s \\ y &= q_{\star} + \beta_1 x + \beta_2 y + \beta_3 s \\ s &= t_{\star} + \gamma_1 x + \gamma_2 y + \gamma_3 s \end{aligned}$$

where the rotation matrix is the transpose of the rotation matrix that appears in (5), and:

$$(12) \quad \begin{aligned} p_{\star} &= -\alpha_1 x_{\star} - \alpha_2 y_{\star} - \alpha_3 s_{\star} \\ q_{\star} &= -\beta_1 x_{\star} - \beta_2 y_{\star} - \beta_3 s_{\star} \\ t_{\star} &= -\gamma_1 x_{\star} - \gamma_2 y_{\star} - \gamma_3 s_{\star} \end{aligned}$$

To obtain the particle parameters at the input of the following drift space it is now enough to apply the formulae of section 3 changing the vector (x_0, x'_0, y_0, y'_0) with $(\xi_f, \xi'_f, \eta_f, \eta'_f)$ and the transformation (5) with the transformation (5').

Section 9

On each of the two dimensional phase-spaces the beam can be described by the parameters of the enveloping ellipse.

That is, one assumes that in the four-dimensional phase space (x, z, x', z') the whole envelope is the product of two independent functions whose arguments are respectively (x, x') , (z, z') .

Assuming for example, that in the two phase planes, the ellipses are as shown in fig. 5 the beam, in actual space, has a non continuous external envelope, whose intersection with the xz plane is a rectangular surface from every point of which ∞^2 paths, contained within a solid angle with rectangular cross section, start (see fig. 6).

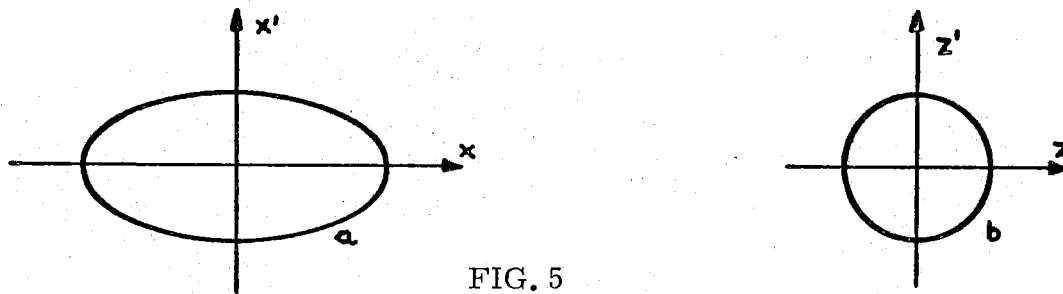


FIG. 5

Such a solid angle vanishes for the outer particles.

For a clear understanding of how the situation develops from one section to the next let us first of all consider the trajectory of each particle.

As long as the channel is errorless, the final position of a particle in the xx' plane only depends upon the projection, onto this plane, of its initial position; the same is true for the vertical plane. All this still holds in presence of translation errors, or length variations of drift spaces or energy variations. But when there are errors due to rotations, the two phase planes are no longer independent.

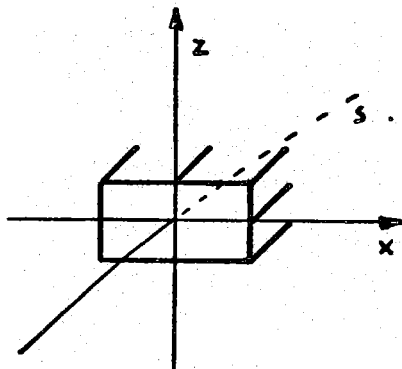


FIG. 6

That is to say (see section 3) that the trajectory of a particle depends upon all four initial coordinates, in both phase planes.

Considering again the elliptic envelopes, as a consequence of what we said before, it is clear that:

- 1) - the developments of the initial ellipses (in each of the two phase planes) depend upon the initial conditions that one assumes on the other plane.

- 2) - Since the transformations of section 3 are non linear, an ellipse does not remain one as a geometrical locus.

Referring to point 1) we note that initial points, internal to the envelope, are transformed into internal ones and external points into external ones in the whole phase space (since the transformations between initial and final points are continuous). Therefore in order to see how much the final envelope is changed, it is sufficient to study only points (of the initial envelope) that satisfy the equations:

$$f(x, x') = 0$$

$$g(z, z') = 0$$

Considering only envelope points, the initial conditions to be examined reduce from ∞^4 to ∞^3 . Besides such conditions reduce again to ∞^2 since the initial envelope is the product of the two independent functions f and g .

Given the philosophy of our computation it is useful to group the initial conditions into ∞^1 ellipses; i. e. the ellipse $f(x, x') = 0$ together with the ∞^1 points which satisfy the equation $g(z, z') = 0$ and, viceversa, the ellipse $g(z, z') = 0$ together with the ∞^1 points which satisfy the equations $f(x, x') = 0$.

We furthermore note that the effect mentioned under point 2) can be neglected since, a posteriori, the variations of the final ellipse parameters, due to channel errors, give small effects and therefore an accurate quantitative evaluation of the change in the shape of the ellipse is not justified in view of the difficulty of the computation.

However, even though we neglect point 2) the computation of the effects of the errors of a magnetic channel on the final ellipse parameters is very complicated.

We have therefore chosen to start the computation with a finite set of points on the ellipse at the channel input and to then reconstruct the ellipse by means of the transformed points at the channel output.

Section 10

In this section we will describe the practical method we used. For each initial ellipse we chose the 8 intersections, of the ellipse with the axes and with their bisectrices.

In the four dimensional phase space we have considered the 64 points obtained relating the 8-point sets to each other among themselves.

Let us indicate each point with a two figure number, the first figure indicating the point (in two dimension) of the first ellypse (see fig. 7) and the second one the point of the second ellypse.

At the channel output we have considered the following 8-point transformed set:

11	21	31	81
12	22	32	82
.....				
18	28	38	88

to obtain the ellypses of the horizontal plane, and the following 8-point set:

11	12	13	18
21	22	23	28
.....				
81	82	83	88

to obtain the ellypses of the vertical plane.

Since we, (according to what said before), neglect the change of shape of the ellypse, 8 points are sufficient to describe such a geometrical locus. We have used such points in the following manner (see fig. 7).

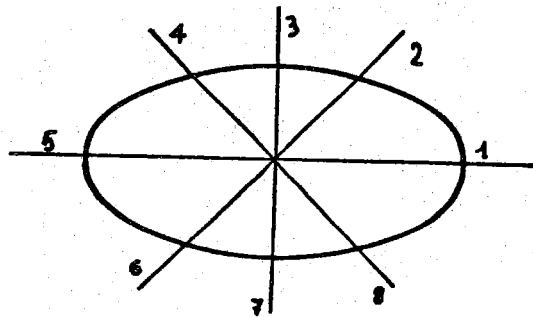


FIG. 7

For each set of 8 points we have defined the center P_0 of the ellypse by the weighted average of points 1, 2, 3, 4, 5, 6, 7, 8, then we calculated the centered ellypse through the points:

$$Q_1 = \frac{\vec{P}_1 - \vec{P}_5}{2} \qquad Q_2 = \frac{\vec{P}_2 - \vec{P}_6}{2} \qquad Q_3 = \frac{\vec{P}_3 - \vec{P}_7}{2}$$

Now to obtain the equation $f(x, x')$ of the ellipse it is sufficient to compute the determinant:

$$\begin{vmatrix} x^2 & x x' & x'^2 & 1 \\ x_1^2 & x_1 x'_1 & x_1'^2 & 1 \\ x_2^2 & x_2 x'_2 & x_2'^2 & 1 \\ x_3^2 & x_3 x'_3 & x_3'^2 & 1 \end{vmatrix} = 0$$

To compute the effects of the channel errors we have used the following procedure.

The channel has been first traced with errorless elements, computing the trajectory of the 8 points.

At the channel output, (as we said before), we define:

- a) - the center coordinates of the ellipse
- b) - the R and X parameters
- c) - maximum half-displacement and maximum half-inclination.

For each kind of element error the channel is traced following the 64-point paths and the parameters a), b), c) are redefined as said before. Such parameters will present small variations compared with the ones obtained for the errorless channel. This enables us to compute the incremental ratios that we will consider as derivatives.

Let us now remark that the computed derivatives should also be used to shift the beam at the channel output as much as one wishes.

Section 11

In this section we will give some remarks about the subroutines used by the program.

- 1) - DM
- 2) - QM
- 3) - MH
- 4) - MV

The subprograms transform the vectors at the input face into vectors at the output face applying the element matrices of drift spaces, quadrupoles and bending magnet.

18.

They are used in the block diagram at points marked 1, 2, 3, 4, 5, 6.

5) - SILVA

The subprogram calculates eight bi-dimensional vectors by means of the E, R, X ellipse parameters as said in section 9. It is used, in the block diagram at point 7 both for the initial horizontal ellipse and for the vertical one.

6) - ELLI

The subroutine computes the ellipse parameters (Y, Y', R, X) from a set of eight points using section 10 formulae.

This program is used at point 8 of the block diagram.

7) - FORVE

This subprogram, making use of the two sets of eight bi-dimensional vectors produces 64 quadri-dimensional vectors as said before (section 10). This subprogram is also used at point 7 of the block diagram. All of the following subroutines are used at points 5 and 6.

8) - DEF

This subprogram, according KEK values, defines the three translation (KEK=4, 5, 6) and rotation (KEK=1, 2, 3) parameters of the displaced element (as said in section 4 for the \mathcal{Z} transformation). On the contrary for KEK=7 this subroutine introduces the quantity $\Delta i/i \neq 0$ for the purposes have mentioned in section 2.

9) - COOR

The subprogram computes the initial and final coordinates referred to \overline{T}_1 and T_3 for bending magnets and quadrupoles. According to what said in the second part of section 3.

10) - CALCO

This subprogram computes the parameters of coordinate transformation \mathcal{F} of $T_1(T_3)$ into $\overline{T}_1(\overline{T}_3)$ by means of section 4 formulae.

11) - SPOST

Such a subroutine computes the vectors at the input (output) of displaced element making use of section 6 formulae.

12) - RICO

Given the transformation that transforms a system A into a system B, the subprogram RICO computes the transformation that transforms the system B into A, according to formulae 5'), 12) of section 8.

Section 12

We will now present a sample program output.

The example refers to the transport channel which transfers the electrons (positrons) from the Linac to the machine "Adone". The channel consists of 7 bending magnets, 14 quadrupole lenses, and 20 drift spaces.

The elements constituting the channel are first described row by row as they are geometrically placed into the channel.

On each row the following parameters are written:

- a) - for magnets:
field index, radius of curvature, tangent of the slope at the input face, and at the output face, geometrical deflection angle (degrees)
- b) - quadrupoles:
strength (which in thin lens approximation coincides with the a_{21} matrix element), length of quadrupole
- c) - for drift spaces:
length of the drift space.

The final errorless parameters of the beam are typed next. Another set of rows is written after this.

In each row the values of the derivatives of the two center coordinates, the R, X quantities and the two maximum displacements, with respect to varied parameter are typed.

We have defined the latter six quantities in section 1.

On the right of each row there are:

- a) - the order number of the channel element
- b) - the type of variation:
 - RX, RY, RZ rotations
 - TX, TY, TZ translations
 - VE variation of the current supply
 - VL length variation.
- c) - the channel element type:
 - MG bending magnet
 - TD drift space
 - QP quadrupole
- d) - the phase plane
 - O horizontal
 - V vertical.

Section 13

By now analyzing the numerical results we deduce that:

- 1) - the variations of the dimensions of the final ellipses are negligible; this may be deduced by comparing the first and 5th column with the second and 6th.

One sees that either the differences are small or, in case the results of the first and the second column vanish, the results of the 5th and 6th columns are negligible in absolute value.

This may be interpreted by saying that the displacement of one element is equivalent to the addition (subtraction) at the element input (at the element output) of a vector which has little dependence from the input conditions.

- 2) - However, in case one wants, to take into account the ellipse deformations, we list below the cases for which the deformation is more relevant.

For the quadrupoles RZ, VE and TZ are effective (notice that VE and TZ do not shift the ellipse center).

For the magnets RZ and TX are irrelevant.

- 3) - About the center shift:

For the quadrupoles only TX and TY produce relevant effects. TZ and VE, are indeed ineffective, RZ has only a small effect (though non zero, as we are in phase-space), while finally RX and RY give opposite effects at the element extremes.

For the magnets TX and TY whose effects however, are smaller than for the quadrupoles (because of the low field indices are important. VE and RZ have little importance.

From the data reported below one can observe for instance that a 1 mm x-displacement of element 37 (quadrupole with $a_{21} \simeq -1$) produces a horizontal displacement of about 5 mm of the ellipse center.

On the other hand a similar Y-displacement of element 21 (quadrupole with $a_{21} \simeq 0.77$) shifts the ellipse center vertically of about 9,35 mm.

As far as the slopes are conserved we call attention on element 23 (quadrupole with $a_{21} \simeq 0.74$) that, with a 1 mm x-translation, shifts the center of 2.19 mrad, and on element 21 that with a 1 mm y-translation shifts the center of 3.2 mrad.

The effects we have reported are the largest in absolute value.

- 4) - One element is particularly sensitive to the VL variation (the 36th drift space 70 cm long).

In this case a 1 cm variation produces variations in the ellipse dimensions of 0.24 and 0.1 mrad (horizontally) and 0.4 and 0.23 mrad (vertically).

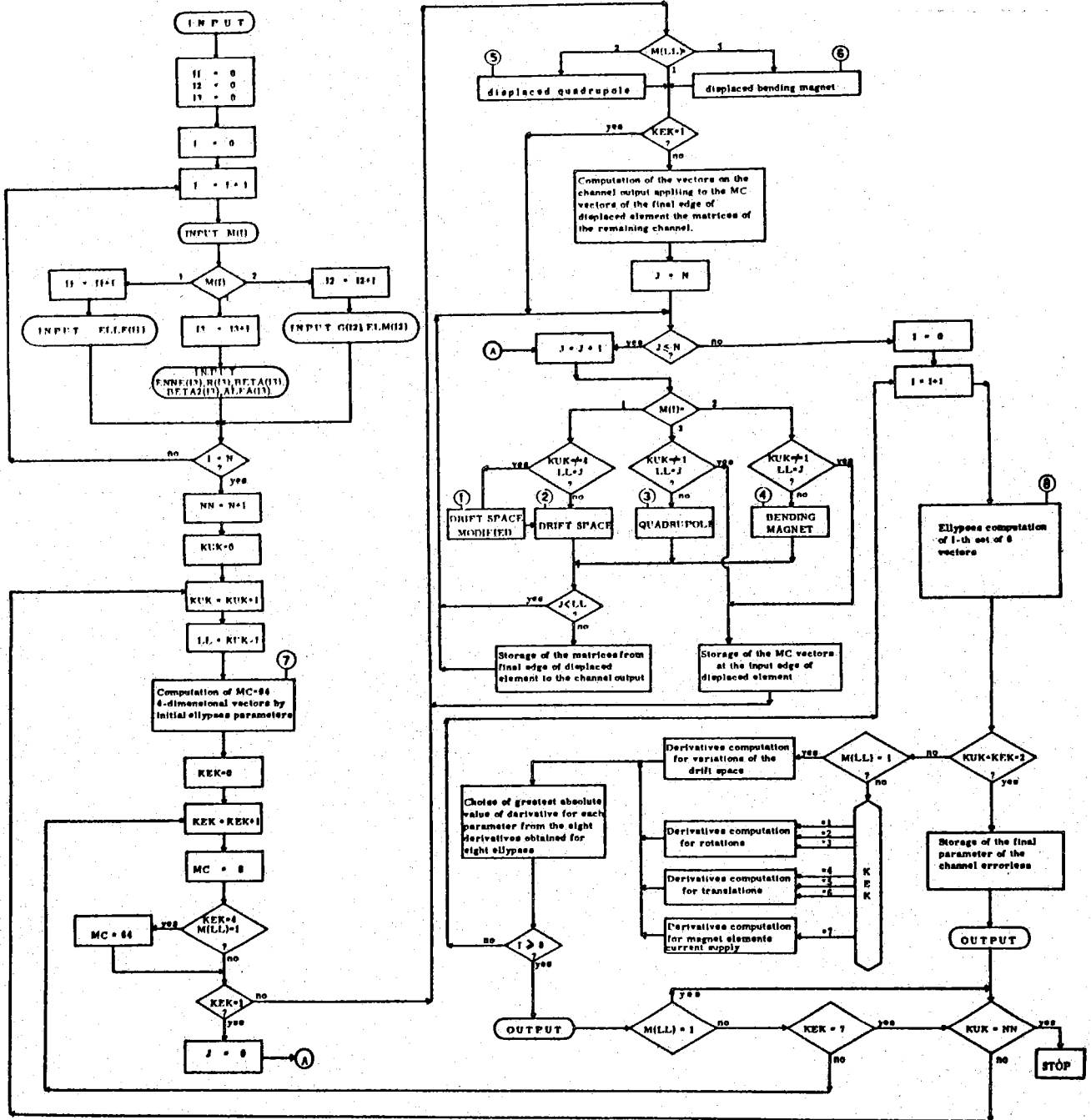
In conclusion one must be careful with translations perpendicular to the channel for quadrupoles and for bending magnets.

For magnets there are also small variations of RZ in horizontal and VE in vertical that produce relevant effects.

The authors are much indebted to F. Amman for his helpful criticism.

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- (2) - J. Fronteau: "Le théorème de Liouville et le problème général de la stabilité" CERN 65/38 (1965).
- (3) - N. D. Sen Gubta: "On the motion of a charged particle in a magnetic field" *Nuovo Cimento* 42, (1966).



CHANNEL COMPOSITION

```

*****
1  MAGNET          ..... 0.          0.51110D 01 0.          0.10500D 00 0.60000D 01
2  DRIFT SPACE    ..... 0.25383D 01          0.12000D 00
3  QUADRUPOLE     ..... 0.29928D 00
4  DRIFT SPACE    ..... 0.35000D 00
5  MAGNET          ..... -0.17000D 01 0.12710D 01 0.          0.          0.24000D 02
6  DRIFT SPACE    ..... 0.40000D 00
7  QUADRUPOLE     ..... 0.53310D 00 0.21000D 00
8  DRIFT SPACE    ..... 0.24557D 01
9  QUADRUPOLE     ..... -0.63701D 00 0.32000D 00
10 DRIFT SPACE    ..... 0.24557D 01
11 QUADRUPOLE     ..... 0.47860D 00 0.21000D 00
12 DRIFT SPACE    ..... 0.40000D 00
13 MAGNET          ..... -0.17000D 01 0.14020D 01 0.          0.          0.30000D 02
14 DRIFT SPACE    ..... 0.40000D 00
15 QUADRUPOLE     ..... 0.50435D 00 0.21000D 00
16 DRIFT SPACE    ..... 0.12750D 01
17 QUADRUPOLE     ..... 0.16880D 00 0.12000D 00
18 DRIFT SPACE    ..... 0.54500D 01
19 QUADRUPOLE     ..... -0.25084D 00 0.12000D 00
20 DRIFT SPACE    ..... 0.39500D 01
21 QUADRUPOLE     ..... 0.77035D 00 0.32000D 00
22 DRIFT SPACE    ..... 0.75000D 00
23 QUADRUPOLE     ..... -0.74101D 00 0.32000D 00
24 DRIFT SPACE    ..... 0.22201D 01
25 MAGNET          ..... 0.          0.12710D 01 0.          0.          0.26250D 02
26 DRIFT SPACE    ..... 0.35000D 00
27 QUADRUPOLE     ..... 0.25000D 00 0.12000D 00
28 DRIFT SPACE    ..... 0.16434D 01
29 QUADRUPOLE     ..... -0.95435D 00 0.39000D 00
30 DRIFT SPACE    ..... 0.16434D 01
31 QUADRUPOLE     ..... 0.19486D 00 0.12000D 00
32 DRIFT SPACE    ..... 0.35000D 00
33 MAGNET          ..... 0.          0.14630D 01 0.          0.          0.26250D 02
34 DRIFT SPACE    ..... 0.17967D 01
35 QUADRUPOLE     ..... 0.86865D 00 0.39000D 00
36 DRIFT SPACE    ..... 0.70000D 00
37 QUADRUPOLE     ..... -0.95800D 00 0.39000D 00
38 DRIFT SPACE    ..... 0.40000D 01
39 MAGNET          ..... 0.          0.13751D 02 0.          0.          0.37500D 01
40 DRIFT SPACE    ..... 0.70000D-01
41 MAGNET          ..... 0.          0.13751D 02 0.          0.          0.37500D 01

```


XC	XCP	R	X	XMAX	XPHAX				
0.	0.	0.86953D 00	0.65385D 00	0.36894D-02	0.29488D-02	0			
0.	0.	0.14760D 01	0.90035D 00	0.45003D-02	0.38419D-02	V			
DXC	DXCP	DR	DX	DXMAX	DXPHAX	N			
-0.33548D-03	0.33923D-03	-0.12069D-03	-0.70305D-04	0.33508D-03	0.33981D-03	1	RX	MG	0
-0.27066D-02	-0.13111D-01	-0.53375D-05	0.39276D-05	0.27066D-02	0.13111D-01	1	RX	MG	V
-0.13200D-01	0.39489D-02	0.20772D-05	0.59061D-05	0.13200D-01	0.39489D-02	1	RY	MG	0
0.20329D-14	-0.21684D-14	0.56031D-03	0.86559D-03	0.15642D-05	0.72921D-06	1	RY	MG	V
-0.52152D-03	0.12200D-03	0.67052D-04	0.83901D-04	0.52173D-03	0.12211D-03	1	RZ	MG	0
-0.51058D-01	-0.24746D 00	-0.11352D-04	0.74776D-05	0.51058D-01	0.24746D 00	1	RZ	MG	V
0.15391D-02	0.18191D-02	-0.44577D-02	-0.79926D-02	0.15202D-02	0.18115D-02	1	TX	MG	0
0.11520D-15	-0.16263D-15	-0.13117D-01	0.78388D-02	0.14737D-05	-0.17071D-04	1	TX	MG	V
0.43368D-15	0.47434D-15	-0.45253D-08	-0.18427D-07	-0.40221D-10	-0.76722D-11	1	TY	MG	0
-0.75493D-02	-0.47837D-01	0.28586D-07	-0.54905D-07	0.75493D-02	0.47837D-01	1	TY	MG	V
0.30313D-01	0.34449D-01	0.43555D-03	0.85323D-03	0.30315D-01	0.34450D-01	1	TZ	MG	0
0.11520D-15	0.	0.33693D-01	0.50416D-01	0.91842D-04	0.43849D-04	1	TZ	MG	V
-0.24804D 00	0.80048D-01	-0.45242D-09	-0.18424D-08	0.24804D 00	0.80048D-01	1	VE	MG	0
0.19651D-16	-0.10842D-16	0.28587D-08	-0.54904D-08	-0.54475D-11	0.37205D-11	1	VE	MG	V
0.54210D-17	-0.81315D-17	-0.42837D-01	-0.76756D-01	-0.18167D-03	-0.72639D-04	2	VL	TD	0
0.10842D-16	0.	-0.14214D 00	0.49660D-01	-0.31789D-04	-0.18501D-03	2	VL	TD	V
0.54210D-15	-0.16263D-14	-0.91531D-02	0.18287D-01	0.26686D-04	-0.19672D-04	3	RX	QP	0
0.12505D-03	0.16135D-03	-0.10797D-03	0.13437D-04	0.12499D-03	0.16121D-03	3	RX	QP	V
-0.27551D-03	-0.13003D-03	-0.65810D-05	0.70659D-04	0.27565D-03	0.13002D-03	3	RY	QP	0
0.24395D-14	0.10842D-14	-0.22688D-01	0.16089D-01	-0.27331D-05	-0.35428D-04	3	RY	QP	V
0.18008D-01	-0.92974D-03	-0.26201D-04	0.28092D-03	0.18009D-01	0.92970D-03	3	RZ	QP	0
0.46652D-02	-0.11560D-01	-0.43452D-03	0.54045D-04	0.46649D-02	0.11559D-01	3	RZ	QP	V
0.95572D 00	0.49362D-01	-0.18763D-10	-0.19762D-10	0.95572D 00	0.49362D-01	3	TX	QP	0
0.18974D-15	0.54210D-16	-0.58842D-10	-0.51403D-10	-0.11102D-12	-0.76328D-13	3	TX	QP	V
0.	-0.54210D-16	-0.20872D-10	-0.23315D-10	-0.59848D-13	-0.35562D-13	3	TY	QP	0
0.23314D 00	-0.57778D 00	-0.58398D-10	-0.51847D-10	0.23314D 00	0.57778D 00	3	TY	QP	V
0.	-0.67763D-16	0.61583D 00	-0.96418D 00	-0.16011D-02	0.10441D-02	3	TZ	QP	0
0.10164D-15	0.	-0.10839D 01	0.10191D 01	0.62613D-03	-0.14108D-02	3	TZ	QP	V
0.	-0.20329D-16	-0.26336D 00	0.28269D 01	0.56562D-02	-0.44673D-03	3	VE	QP	0
0.10164D-16	0.16263D-16	-0.42455D 01	0.49307D 00	-0.22602D-02	-0.55452D-02	3	VE	QP	V
-0.54210D-17	-0.54210D-17	-0.65667D 00	0.88597D 00	0.14325D-02	-0.11145D-02	4	VL	TD	0
0.54210D-17	0.	0.93920D 00	-0.97536D 00	-0.65803D-03	0.12213D-02	4	VL	TD	V
-0.21472D-02	-0.81028D-03	0.10098D-01	-0.26559D-01	0.21210D-02	0.83382D-03	5	RX	MG	0
-0.13073D-01	-0.65288D-02	0.10339D-03	-0.34050D-06	0.13073D-01	0.65289D-02	5	RX	MG	V
-0.98089D-02	0.55734D-02	-0.23089D-01	-0.90565D-01	0.96107D-02	0.55342D-02	5	RY	MG	0
0.10842D-14	-0.21684D-14	0.20245D-01	-0.82637D-01	-0.79594D-04	0.38320D-04	5	RY	MG	V
-0.41567D-01	-0.42160D-02	0.20358D-02	-0.13363D-01	0.41589D-01	0.42126D-02	5	RZ	MG	0
0.41082D 00	-0.66990D 00	0.68535D-03	0.67381D-05	0.41082D 00	0.66990D 00	5	RZ	MG	V
-0.31777D 01	-0.32279D 00	-0.10057D 00	0.10448D 01	0.31797D 01	0.32262D 00	5	TX	MG	0
0.54210D-16	-0.10842D-15	0.27884D 00	-0.44919D-01	0.13367D-03	0.36288D-03	5	TX	MG	V
-0.23419D-13	0.13146D-13	-0.95912D-07	0.30684D-06	0.56892D-09	-0.16262D-09	5	TY	MG	0
-0.54032D 00	0.88097D 00	0.40285D-06	-0.33904D-06	0.54032D 00	0.88097D 00	5	TY	MG	V
-0.15741D 00	0.89150D-01	-0.81476D 00	-0.31798D 01	0.15045D 00	0.87769D-01	5	TZ	MG	0
0.	0.54210D-16	0.52982D 00	-0.16869D 01	-0.19156D-02	0.68949D-03	5	TZ	MG	V
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0.	0.54210D-17	0.40284D-07	-0.33904D-07	-0.17849D-10	0.52427D-10	5	VE	MG	V

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0. 0.176180-15	-0.542100-16 0.108420-15	-0.211320 00 -0.811010 00	0.359650 01 0.141550 01	0.721330-02 0.135410-02	-0.358330-03 -0.105550-02	7 7	TZ TZ	QP QP	D V
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0.518740-02 0.675160-02	-0.512250-02 0.373220-03	0.509020-04 -0.224530-04	0.510980-04 0.121000-03	0.518750-02 0.675180-02	0.512260-02 0.373190-03	9 9	RZ RZ	QP QP	D V
-0.709740 00 0.542100-16	-0.700230 00 -0.542100-16	-0.189850-10 0.119900-10	-0.182080-10 -0.499600-11	0.709740 00 0.173470-14	0.700230 00 0.160460-13	9 9	TX TX	QP QP	D V
0. -0.964070 00	-0.271050-16 -0.546070-01	-0.187630-10 0.137670-10	-0.194290-10 -0.544010-11	-0.503070-13 0.964070 00	-0.320920-13 0.546070-01	9 9	TY TY	QP QP	D V
0.108420-15 0.108420-15	-0.813150-16 0.108420-15	0.112730 01 0.189990 01	-0.386230 00 0.428800 00	-0.122140-03 0.190670-02	0.191110-02 0.247220-02	9 9	TZ TZ	QP QP	D V
0.108420-16 0.	-0.135530-16 -0.542100-17	0.478880 00 -0.228650 00	0.477750 00 0.129110 01	0.125480-02 0.159670-02	0.811450-03 -0.297630-03	9 9	VE VE	QP QP	D V
-0.542100-17 0.108420-16	-0.813150-17 0.	-0.738840 00 -0.693020 00	0.925930 00 -0.108600 01	0.146780-02 -0.195590-02	-0.125410-02 -0.902450-03	10 10	VL VL	TD TD	D V
-0.108420-14 0.124440-03	-0.135530-14 0.526400-03	0.678680-02 0.142590-03	0.402500-01 0.370430-04	0.707260-04 0.124590-03	0.173580-04 0.526590-03	11 11	RX RX	QP QP	D V
0.312800-02 0.	0.674650-03 -0.216840-14	0.532660-05 0.234670-01	0.545890-04 0.930880-02	0.312820-02 0.815830-05	0.674660-03 0.239020-04	11 11	RY RY	QP QP	D V
-0.165590-01 0.260810-01	-0.965290-03 -0.127800-01	0.206240-04 0.579220-03	0.214400-03 0.150780-03	0.165590-01 0.260820-01	0.965320-03 0.127800-01	11 11	RZ RZ	QP QP	D V
-0.992750 00 0.	0.578840-01 -0.542100-16	-0.193180-10 0.106580-10	-0.266450-10 -0.666130-11	0.992750 00 -0.173470-14	0.578840-01 0.143110-13	11 11	TX TX	QP QP	D V
0. 0.182200 01	-0.813150-16 0.892460 00	-0.180970-10 0.754950-11	-0.259790-10 -0.655030-11	-0.641850-13 0.182200 01	-0.303580-13 0.892460 00	11 11	TY TY	QP QP	D V
0. 0.	-0.948680-16 -0.542100-16	0.535850 00 -0.371630 00	-0.321620 01 -0.153190 01	-0.623200-02 -0.233530-02	0.908530-03 -0.483670-03	11 11	TZ TZ	QP QP	D V
0. 0.	-0.813150-17 -0.108420-16	0.214100 00 0.561310 01	0.220800 01 0.153470 01	0.464440-02 0.598950-02	0.362910-03 0.727070-02	11 11	VE VE	QP QP	D V

0.	-0.813150-17	-0.126970 01	0.412810 01	0.782730-02	-0.215690-02	12	VL	TD	0
0.162630-16	0.	-0.317940 00	0.441940 00	0.378650-03	-0.413890-03	12	VL	TD	V
0.213570-02	0.982660-03	0.159110-01	-0.189590-01	0.211360-02	0.101640-02	13	RX	MG	0
-0.129530-01	-0.214750-01	-0.176300-04	0.289920-04	0.129530-01	0.214750-01	13	RX	MG	V
-0.332610-01	-0.162840-01	-0.832960-01	0.122260 00	0.334610-01	0.161430-01	13	RY	MG	0
0.	0.	0.609470-01	0.182770 00	0.314730-03	0.100220-03	13	RY	MG	V
0.411230-01	0.175040-02	0.188410-02	-0.187390-01	0.411560-01	0.174790-02	13	RZ	MG	0
0.183220 01	0.100780 01	-0.508810-03	-0.212720-03	0.183220 01	0.100780 01	13	RZ	MG	V
0.330590 01	0.139840 00	-0.209430 00	0.148010 01	0.330880 01	0.139490 00	13	TX	MG	0
0.	0.	-0.422820 00	-0.482180 00	-0.948570-03	-0.550290-03	13	TX	MG	V
-0.100720-12	-0.493040-13	-0.696600-06	-0.285930-05	-0.623780-08	-0.118110-08	13	TY	MG	0
-0.221880 01	-0.122040 01	0.722420-05	0.755970-05	0.221880 01	0.122040 01	13	TY	MG	V
-0.300490 00	-0.147120 00	-0.168370 01	0.247660 01	0.304560 00	0.144270 00	13	TZ	MG	0
0.	0.	0.960800 00	0.265530 01	0.427000-02	0.125030-02	13	TZ	MG	V
0.177710 01	0.751890-01	-0.696610-07	-0.285930-06	0.177710 01	0.751890-01	13	VE	MG	0
0.	0.	0.722420-06	0.755970-06	0.152870-08	0.940170-09	13	VE	MG	V
-0.542100-17	-0.271050-17	0.468270 00	0.157830 01	0.349610-02	0.793480-03	14	VL	TD	0
0.542100-17	0.	-0.133160 01	-0.228940 01	-0.403280-02	-0.173490-02	14	VL	TD	V
-0.542100-15	-0.230390-14	-0.718880-02	0.255510-01	0.647870-04	-0.189830-04	15	RX	QP	0
0.594750-03	0.805430-03	0.184330-03	0.133260-03	0.595060-03	0.805670-03	15	RX	QP	V
-0.680450-03	0.548770-03	-0.201530-04	0.953910-04	0.680630-03	0.548730-03	15	RY	QP	0
0.	0.108420-14	0.283640-01	0.444890-01	0.901710-04	0.455200-04	15	RY	QP	V
-0.277200-01	-0.307680-02	-0.795280-04	0.374920-03	0.277210-01	0.309670-02	15	RZ	QP	0
0.377630-01	-0.228260-01	0.749600-03	0.542620-03	0.377640-01	0.228270-01	15	RZ	QP	V
-0.151850 01	-0.169600 00	-0.175420-10	-0.190960-10	0.151850 01	0.169600 00	15	TX	QP	0
0.	0.	0.111020-11	-0.366370-11	-0.346940-14	0.173470-14	15	TX	QP	V
-0.542100-16	-0.189740-15	-0.166530-10	-0.164310-10	-0.433680-13	-0.277560-13	15	TY	QP	0
0.211920 01	0.128110 01	-0.399680-11	0.144330-11	0.211920 01	0.128110 01	15	TY	QP	V
0.542100-16	-0.813150-16	0.559220 00	0.166740 01	0.372800-02	0.948140-03	15	TZ	QP	0
0.	0.	-0.970320 00	-0.250500 01	-0.407190-02	-0.126290-02	15	TZ	QP	V
0.	-0.176180-16	-0.815590 00	0.385730 01	0.749990-02	-0.138450-02	15	VE	QP	0
0.	-0.108420-16	0.716390 01	0.545280 01	0.123090-01	0.926750-02	15	VE	QP	V
0.	-0.677630-17	-0.919950-01	-0.906580-01	-0.238990-03	-0.156010-03	16	VL	TD	0
0.108420-16	0.	-0.346320 00	0.200650 00	0.309880-04	-0.450840-03	16	VL	TD	V
0.108420-14	-0.230390-14	-0.681330-02	0.610300-02	0.646450-05	-0.132510-04	17	RX	QP	0
-0.208100-03	-0.812950-04	0.350860-04	0.424200-04	0.208190-03	0.813410-04	17	RX	QP	V
0.137230-03	0.928520-04	-0.175450-04	0.339440-04	0.137290-03	0.928220-04	17	RY	QP	0
0.	-0.108420-14	0.149400-01	0.726200-02	0.168100-04	0.164930-04	17	RY	QP	V
-0.837960-02	-0.198080-02	-0.699540-04	0.135310-03	0.837980-02	0.198070-02	17	RZ	QP	0
0.100930-01	0.700490-02	0.143990-03	0.170250-03	0.100930-01	0.700510-02	17	RZ	QP	V
-0.584780 00	-0.138250 00	-0.156540-10	-0.263120-10	0.584780 00	0.138250 00	17	TX	QP	0
0.	0.	0.304200-10	-0.300870-10	-0.199490-13	0.398990-13	17	TX	QP	V
0.	-0.162630-15	-0.148770-10	-0.268670-10	-0.641850-13	-0.251530-13	17	TY	QP	0
0.560460 00	0.388990 00	0.488500-10	-0.490720-10	0.560460 00	0.388990 00	17	TY	QP	V
0.	-0.542100-16	0.480200 00	-0.287350 00	-0.302550-03	0.814100-03	17	TZ	QP	0
0.	-0.108420-15	0.100190 01	0.418240 00	0.126720-02	0.130640-02	17	TZ	QP	V
0.	-0.108420-16	-0.698640 00	0.135310 01	0.237030-02	-0.118580-02	17	VE	QP	0
0.	0.	0.141550 01	0.169610 01	0.328310-02	0.184000-02	17	VE	QP	V

0. 0.10842D-16	-0.81315D-17 0.	-0.57089D 00 -0.13480D 01	0.19695D 00 -0.21316D 00	0.68270D-04 -0.12288D-02	-0.96880D-03 -0.17563D-02	18 18	VL VL	TD TD	0 V
-0.10842D-14 0.53370D-03	-0.18974D-14 0.22160D-03	0.40642D-01 -0.52328D-04	-0.37578D-01 0.39372D-04	-0.56432D-04 0.53372D-03	0.65894D-04 0.22153D-03	19 19	RX RX	QP QP	0 V
0.39551D-04 0.	-0.25747D-04 0.	0.25213D-03 0.65778D-01	-0.19032D-03 0.24180D-01	0.39312D-04 -0.23426D-04	0.26174D-04 -0.94347D-04	19 19	RY RY	QP QP	0 V
-0.36799D-01 0.34336D-01	-0.15335D-01 0.88986D-02	0.10135D-02 -0.20828D-03	-0.76506D-03 0.15669D-03	0.36798D-01 0.34337D-01	0.15337D-01 0.88984D-02	19 19	RZ RZ	QP QP	0 V
0.21850D 01 0.	0.91054D 00 0.	-0.32419D-10 0.17542D-10	-0.54068D-10 -0.16986D-10	0.21850D 01 -0.10408D-13	0.91054D 00 0.22985D-13	19 19	TX TX	QP QP	0 V
0. 0.94833D 00	0. 0.24577D 00	-0.19096D-10 -0.68834D-11	-0.29976D-10 0.87708D-11	-0.72858D-13 0.94833D 00	-0.32092D-13 0.24577D 00	19 19	TY TY	QP QP	0 V
0. 0.	0. 0.	0.78278D-01 0.23738D 01	-0.60293D 00 -0.85371D 00	-0.11825D-02 0.50095D-03	0.13273D-03 0.30887D-02	19 19	TZ TZ	QP QP	0 V
0. 0.	0. 0.	0.10392D 02 -0.20830D 01	-0.76068D 01 0.15729D 01	-0.76750D-02 0.70912D-03	0.17365D-01 -0.27157D-02	19 19	VE VE	QP QP	0 V
-0.54210D-17 0.10842D-16	-0.12197D-16 0.	-0.64854D 00 -0.36960D 01	0.79839D 00 0.65227D 00	0.12572D-02 -0.16638D-02	-0.11007D-02 -0.48253D-02	20 20	VL VL	TD TD	0 V
-0.10842D-14 0.16354D-01	-0.20329D-14 0.50302D-02	-0.11827D 00 0.20698D-03	0.18585D-01 -0.66229D-03	0.14556D-03 0.16354D-01	-0.23112D-03 0.50329D-02	21 21	RX RX	QP QP	0 V
0.24349D-02 0.	0.29198D-02 0.	-0.19006D-03 0.53606D 00	0.67530D-04 -0.20276D 00	0.24349D-02 0.83048D-04	0.29195D-02 0.68387D-03	21 21	RY RY	QP QP	0 V
0.17125D 00 -0.17284D 00	0.90794D-01 -0.59147D-01	-0.72995D-03 0.85979D-02	0.25904D-03 -0.27514D-02	0.17125D 00 0.17284D 00	0.90793D-01 0.59158D-01	21 21	RZ RZ	QP QP	0 V
-0.31856D 01 0.	-0.16893D 01 0.	-0.26867D-10 0.14122D-09	-0.33640D-10 -0.10192D-09	0.31856D 01 -0.39899D-13	0.16893D 01 0.18388D-12	21 21	TX TX	QP QP	0 V
0. -0.93538D 01	0. -0.32009D 01	-0.27089D-10 0.70388D-10	-0.30753D-10 -0.30753D-10	-0.78930D-13 0.93538D 01	-0.45970D-13 0.32009D 01	21 21	TY TY	QP QP	0 V
0. 0.	0. 0.	0.18887D 01 0.31109D 02	0.81151D 00 -0.12130D 02	0.27660D-02 0.56213D-02	0.32017D-02 0.40380D-01	21 21	TZ TZ	QP QP	0 V
0. 0.	0. 0.	-0.76761D 01 0.93362D 02	0.28637D 01 -0.18664D 02	0.19218D-02 0.54533D-01	-0.13163D-01 0.11317D 00	21 21	VE VE	QP QP	0 V
0. 0.16263D-16	-0.94868D-17 0.	-0.25207D 01 -0.32189D 02	0.83328D-02 0.13344D 02	-0.14402D-02 0.22342D-03	-0.42897D-02 -0.43101D-01	22 22	VL VL	TD TD	0 V
-0.54210D-15 -0.26442D-01	-0.20329D-14 -0.83990D-02	0.13647D 00 -0.72993D-03	-0.28782D-02 0.20785D-03	0.98225D-04 0.26442D-01	0.24989D-03 0.83980D-02	23 23	RX RX	QP QP	0 V
-0.20398D-03 0.	-0.13197D-02 0.	0.32751D-03 -0.62556D 00	-0.82677D-04 0.22127D 00	0.20400D-03 -0.11634D-03	0.13202D-02 -0.79692D-03	23 23	RY RY	QP QP	0 V
-0.13085D 00 0.13623D 00	-0.73902D-01 0.47720D-01	0.13589D-02 -0.28158D-02	-0.34312D-03 0.80162D-03	0.13085D 00 0.13623D 00	0.73904D-01 0.47717D-01	23 23	RZ RZ	QP QP	0 V
0.38810D 01 0.	0.21926D 01 0.	0.20539D-10 -0.14211D-10	-0.20095D-10 0.37748D-11	0.38810D 01 -0.43368D-14	0.21926D 01 -0.18648D-13	23 23	TX TX	QP QP	0 V
0. 0.56457D 01	0. 0.19776D 01	0.13767D-10 -0.21538D-10	-0.32196D-10 0.10880D-10	-0.57246D-13 0.56457D 01	0.23419D-13 0.19776D 01	23 23	TY TY	QP QP	0 V
0. 0.	0. 0.	-0.48154D 00 -0.33145D 02	-0.11554D 01 0.11923D 02	-0.26384D-02 -0.65722D-02	-0.81657D-03 -0.43257D-01	23 23	TZ TZ	QP QP	0 V
0. 0.	0. 0.	0.13087D 02 -0.28542D 02	-0.29000D 01 0.92799D 01	0.29670D-02 -0.43769D-02	0.21788D-01 -0.38089D-01	23 23	VE VE	QP QP	0 V

0.54210D-17 0.16263D-16	-0.54210D-17 0.	-0.20411D 01 -0.12355D 01	0.11458D 01 0.95126D 00	0.11889D-02 0.43828D-03	-0.34711D-02 -0.16096D-02	24 24	VL VL	TD TD	O V
0.64527D-02 -0.21004D-01	0.25918D-02 -0.66710D-02	-0.57265D-02 -0.41415D-04	-0.21967D-02 0.17678D-04	0.64605D-02 0.21004D-01	0.26014D-02 0.66710D-02	25 25	RX RX	MG MG	O V
0.31271D-01 0.	0.16070D-01 0.10842D-14	0.24396D-01 -0.34903D-02	0.11762D-01 0.10319D-02	0.31309D-01 -0.99072D-06	0.16111D-01 -0.28211D-05	25 25	RY RY	MG MG	O V
0.18329D-02 -0.16690D 01	0.23005D-02 -0.74042D 00	0.73565D-02 -0.20870D-03	-0.42966D-02 0.82476D-05	0.18345D-02 0.16690D 01	0.22966D-02 0.74042D 00	25 25	RZ RZ	MG MG	O V
0.21829D 00 0.	0.27999D 00 -0.10842D-15	-0.11008D 01 -0.56266D 00	0.43983D 00 0.43233D 00	0.21854D 00 0.19367D-03	0.27813D 00 -0.73229D-03	25 25	TX TX	MG MG	O V
0.23896D-12 -0.11926D-14	0.12289D-12 -0.43368D-15	-0.26369D-05 -0.35324D-05	-0.66255D-06 0.16586D-05	-0.29030D-08 -0.21635D-09	-0.44711D-08 -0.45972D-08	25 25	TY TY	MG MG	O V
0.91727D 00 0.	0.47142D 00 0.	0.71594D 00 -0.35320D-05	0.34550D 00 0.16582D-05	0.91840D 00 -0.21650D-09	0.47263D 00 -0.45966D-08	25 25	TZ TZ	MG MG	O V
0.28478D 00 0.	0.36536D 00 0.	-0.26369D-06 -0.35324D-06	-0.66255D-07 0.16586D-06	0.28478D 00 -0.21635D-10	0.36536D 00 -0.45972D-09	25 25	VE VE	MG MG	O V
0. 0.16263D-16	-0.12197D-16 0.	-0.27655D 01 -0.12355D 01	0.80145D 00 0.95126D 00	0.70554D-04 0.43828D-03	-0.47081D-02 -0.16096D-02	26 26	VL VL	TD TD	O V
-0.10842D-14 0.54139D-03	-0.17618D-14 0.16227D-03	0.30618D-02 0.37824D-04	0.59145D-02 0.12371D-04	0.74187D-05 0.54144D-03	0.74563D-05 0.16232D-03	27 27	RX RX	QP QP	O V
-0.66540D-03 0.	-0.35171D-03 0.10842D-14	0.50981D-06 0.13428D-01	0.37588D-05 0.23137D-02	0.66541D-03 0.98228D-05	0.35171D-03 0.16246D-04	27 27	RY RY	QP QP	O V
0.21877D-02 0.35323D-02	-0.18959D-03 0.17938D-02	0.20012D-05 0.15196D-03	0.14951D-04 0.49733D-04	0.21878D-02 0.35325D-02	0.18959D-03 0.17940D-02	27 27	RZ RZ	QP QP	O V
0.18431D 00 0.	-0.15697D-01 0.10842D-15	-0.17764D-11 -0.26645D-11	0.44409D-12 -0.16764D-10	0.18431D 00 -0.24286D-13	0.15697D-01 -0.34694D-14	27 27	TX TX	QP QP	O V
0.54210D-16 -0.67983D 00	-0.94868D-16 -0.34508D 00	-0.24425D-11 -0.17319D-10	-0.66613D-12 -0.61062D-11	-0.26021D-14 0.67983D 00	-0.43368D-14 0.34508D 00	27 27	TY TY	QP QP	O V
0. 0.	-0.27105D-16 0.	0.35784D 00 0.19266D 01	-0.59477D 00 -0.14115D 00	-0.10010D-02 0.11534D-02	0.60672D-03 0.25070D-02	27 27	TZ TZ	QP QP	O V
0. 0.	-0.14908D-16 0.	0.20259D-01 0.14994D 01	0.15027D 00 0.49646D 00	0.31827D-03 0.17180D-02	0.34343D-04 0.19489D-02	27 27	VE VE	QP QP	O V
-0.54210D-17 0.16263D-16	-0.27105D-17 0.	-0.31136D 01 -0.31442D 01	0.13915D 01 0.10999D 01	0.11121D-02 -0.66843D-03	-0.53034D-02 -0.41029D-02	28 28	VL VL	TD TD	O V
0.10842D-14 -0.29487D-01	-0.21684D-14 -0.71571D-02	0.49236D-01 0.15128D-04	-0.44641D-01 0.22291D-04	-0.66367D-04 0.29487D-01	0.80843D-04 0.71571D-02	29 29	RX RX	QP QP	O V
0.34528D-02 0.	0.41844D-02 -0.54210D-15	0.24172D-03 0.75716D-02	-0.19037D-03 0.60591D-01	0.34526D-02 0.57875D-04	0.41848D-02 0.25766D-04	29 29	RY RY	QP QP	O V
-0.18453D-01 0.14014D-01	-0.75701D-02 -0.38135D-02	0.10212D-02 0.48203D-04	-0.80699D-03 0.81396D-04	0.18452D-01 0.14014D-01	0.75719D-02 0.38136D-02	29 29	RZ RZ	QP QP	O V
-0.40812D 01 0.	-0.16677D 01 0.	0.62172D-11 0.12434D-10	0.30309D-10 -0.20206D-10	0.40812D 01 -0.18215D-13	0.16677D 01 0.16480D-13	29 29	TX TX	QP QP	O V
0. -0.77992D 00	0. 0.21214D 00	-0.88818D-11 0.13323D-10	0.27756D-11 -0.22071D-10	0. 0.77992D 00	-0.14745D-13 0.21214D 00	29 29	TY TY	QP QP	O V
0. 0.	0. -0.54210D-16	-0.18942D 01 0.52170D 00	-0.48924D 00 -0.44057D 01	-0.21123D-02 -0.56010D-02	-0.32127D-02 0.67893D-03	29 29	TZ TZ	QP QP	O V
0. 0.	0. 0.54210D-17	0.93835D 01 0.59478D 00	-0.72050D 01 0.91574D 00	-0.76953D-02 0.16558D-02	0.15702D-01 0.77367D-03	29 29	VE VE	QP QP	O V

-0.542100-17 0.162630-16	-0.162630-16 0.	-0.124750 01 -0.365270 01	0.186990 01 0.547030 01	0.313170-02 0.507780-02	-0.211920-02 -0.476850-02	30 30	VL VL	TD TD	0 V
0.542100-15 0.559540-03	-0.189740-14 0.109340-03	-0.136840-01 0.444340-04	0.767140-02 -0.116640-03	-0.123070-04 0.559410-03	-0.194290-04 0.109400-03	31 31	RX RX	QP QP	0 V
0.375280-03 0.	0.926170-04 -0.108420-14	-0.849030-05 0.167380-01	0.627830-06 -0.342320-01	0.375270-03 -0.329940-04	0.926030-04 0.232580-04	31 31	RY RY	QP QP	0 V
0.836920-02 -0.863110-02	-0.526810-02 -0.967030-03	-0.338550-04 0.178420-03	0.250650-05 -0.468430-03	0.836920-02 0.863060-02	0.526810-02 0.967260-03	31 31	RZ RZ	QP QP	0 V
0.321150 00 0.	0.202100 00 0.108420-15	-0.106580-10 0.102140-10	-0.244250-11 -0.176530-10	0.321150 00 -0.164800-13	0.202100 00 0.134440-13	31 31	TX TX	QP QP	0 V
0. 0.107680 01	-0.542100-16 0.120670 00	-0.943690-11 0.148770-10	-0.344170-11 -0.224270-10	-0.130100-13 0.107680 01	-0.164800-13 0.120670 00	31 31	TY TY	QP QP	0 V
0. 0.	0. -0.108420-15	-0.603400 00 -0.198100 01	0.384940 00 0.359280 01	0.429750-03 0.349570-02	-0.102320-02 -0.257860-02	31 31	TZ TZ	QP QP	0 V
0. 0.	0. 0.	-0.339070 00 0.176950 01	0.254130-01 -0.463790 01	-0.147290-03 -0.494310-02	-0.575210-03 0.229950-02	31 31	VE VE	QP QP	0 V
0.542100-17 0.108420-16	0. 0.	-0.648670 00 -0.169720 01	0.148760 01 0.190040 01	0.267570-02 0.142330-02	-0.110090-02 -0.221190-02	32 32	VL VL	TD TD	0 V
-0.123760-01 -0.331000-01	-0.463850-02 -0.600440-02	-0.904740-02 0.966520-04	0.194860-02 -0.146530-03	0.123850-01 0.331000-01	0.465380-02 0.600450-02	33 33	RX RX	MG MG	0 V
-0.274740-01 0.	-0.756530-02 0.	0.218390-01 0.403320-02	0.309970-02 -0.620060-02	0.274930-01 -0.386250-05	0.760230-02 0.672430-05	33 33	RY RY	MG MG	0 V
-0.295270-02 0.309730 01	-0.386400-02 0.403290 00	-0.861100-02 -0.778680-04	0.102420-01 0.334320-03	0.297210-02 0.309730 01	0.385710-02 0.403290 00	33 33	RZ RZ	MG MG	0 V
-0.189600 00 0.	-0.243220 00 0.	-0.425220 00 -0.773800 00	0.658270 00 0.864630 00	0.190690 00 0.632730-03	0.242500 00 -0.100710-02	33 33	TX TX	MG MG	0 V
-0.212940-12 0.867360-15	-0.586010-13 0.216840-15	-0.249280-05 0.995410-05	-0.174490-06 -0.119400-04	-0.182350-08 -0.923900-08	-0.422690-08 0.129550-07	33 33	TY TY	MG MG	0 V
-0.700110 00 0.	-0.192810 00 0.	0.557030 00 0.995470-05	0.792240-01 -0.119410-04	0.700600 00 -0.923950-08	0.193760 00 0.129550-07	33 33	TZ TZ	MG MG	0 V
-0.284750 00 0.	-0.365350 00 0.	-0.249280-06 0.995410-06	-0.174480-07 -0.119400-05	0.284750 00 -0.923900-09	0.365350 00 0.129550-08	33 33	VE VE	MG MG	0 V
-0.542100-17 0.108420-16	-0.203290-16 0.	-0.121640 01 -0.169720 01	0.140570 01 0.190040 01	0.218780-02 0.142330-02	-0.206610-02 -0.221190-02	34 34	VL VL	TD TD	0 V
-0.108420-14 -0.292360-01	-0.135530-14 -0.356560-02	-0.163650 00 0.956280-03	0.453520 00 -0.160830-02	0.790450-03 0.292350-01	-0.307500-03 0.356690-02	35 35	RX RX	QP QP	0 V
0.313960-01 0.	0.754880-02 0.	-0.368850-04 0.505200 00	0.158870-03 -0.911340 00	0.313970-01 -0.900020-03	0.754880-02 0.642720-03	35 35	RY RY	QP QP	0 V
-0.122340 00 0.155380 00	-0.147810-01 0.247870-01	-0.141140-03 0.401790-02	0.593230-03 -0.676640-02	0.122340 00 0.155370 00	0.147810-01 0.247920-01	35 35	RZ RZ	QP QP	0 V
-0.261860 01 0.	-0.317370 00 0.	-0.888180-11 -0.266450-11	-0.209830-10 0.810460-11	0.261860 01 0.867360-14	0.317370 00 -0.346940-14	35 35	TX TX	QP QP	0 V
0. 0.881560 01	0. 0.140540 01	0.228710-10 -0.559550-10	0.208720-10 0.572880-10	0.555110-13 0.881560 01	0.390310-13 0.140540 01	35 35	TY TY	QP QP	0 V
0. 0.	-0.542100-16 0.	0.434950 01 0.165830 02	-0.116600 02 -0.326060 02	-0.210600-01 -0.319990-01	0.737050-02 0.215510-01	35 35	TZ TZ	QP QP	0 V
0. 0.	0. 0.	-0.154520 01 0.393830 02	0.660190 01 -0.636710 02	0.129070-01 -0.339290-01	-0.262580-02 0.496500-01	35 35	VE VE	QP QP	0 V

-0.542100-17 0.108420-16	-0.542100-17 0.	-0.543410 01 -0.174220 02	0.120650 02 0.335250 02	0.249120-01 0.401860-01	-0.928730-02 -0.230190-01	36 36	VL VL	TD TD	O V
0.108420-14 0.466630-01	-0.176180-14 0.593800-02	0.189470 00 -0.369670-03	-0.474970 00 0.568930-03	-0.831200-03 0.466640-01	0.341470-03 0.593750-02	37 37	RX RX	QP QP	O V
-0.206010-01 0.	-0.554290-02 0.108420-14	0.173890-03 -0.487500 00	-0.592350-03 0.879610 00	0.206000-01 0.883900-03	0.554330-02 -0.608780-03	37 37	RY RY	QP QP	O V
0.150030 00 -0.183530 00	-0.244650-01 -0.313140-01	0.820860-03 -0.140250-02	-0.252040-02 0.214900-02	0.150020 00 0.183530 00	0.244670-01 0.313120-01	37 37	RZ RZ	QP QP	O V
0.524850 01 0.	0.857260 00 0.108420-15	0.111020-11 -0.643930-11	-0.880180-12 0.677240-11	0.524850 01 0.433680-14	0.857260 00 -0.823990-14	37 37	TX TX	QP QP	O V
0. -0.597430 01	-0.406580-16 -0.101880 01	0.103250-10 -0.128790-10	0.182080-10 0.147660-10	0.425010-13 0.597430 01	0.173470-13 0.101880 01	37 37	TY TY	QP QP	O V
0. 0.	-0.406580-16 0.108420-15	-0.556630 01 -0.101340 02	0.120390 02 0.333340 02	0.214350-01 0.332150-01	-0.944590-02 -0.236370-01	37 37	TZ TZ	QP QP	O V
0. 0.	-0.813150-17 -0.108420-16	0.750890 01 -0.152550 02	-0.227330 02 0.238250 02	-0.368130-01 0.256420-01	0.125980-01 -0.201160-01	37 37	VE VE	QP QP	O V
-0.542100-17 0.542100-17	-0.677630-17 0.	0.164610-01 -0.444090-12	0.994620 00 0.100000 01	0.204160-02 0.135820-02	0.279120-04 -0.563790-15	38 38	VL VL	TD TD	O V
0.292940-03 0.509940-03	0.164080-03 -0.214930-03	-0.316780-04 -0.315820-06	0.252880-04 0.227170-06	0.292920-03 0.509940-03	0.164080-03 0.214930-03	39 39	RX RX	MG MG	O V
0.479410-03 0.108420-14	-0.345230-05 -0.108420-14	0.596790-04 -0.128790-09	-0.537630-04 0.960040-05	0.479330-03 0.130130-07	0.355350-05 -0.169140-12	39 39	RY RY	MG MG	O V
0.236920-04 0.929470-01	0.161130-04 0.654560-01	-0.112760-03 -0.486680-07	0.247660-03 -0.161040-06	0.236970-04 0.929470-01	0.161140-04 0.654560-01	39 39	RZ RZ	MG MG	O V
0.674280-02 0.108420-15	0.473180-02 0.	0.809390-03 -0.577320-11	0.653240-01 0.654380-01	0.687640-02 0.886970-04	0.473310-02 -0.737260-14	39 39	TX TX	MG MG	O V
-0.542100-16 0.162630-15	-0.813150-16 0.	-0.710380-07 -0.910380-11	0.316840-07 0.213780-07	0.227420-10 0.289600-10	-0.120450-09 -0.117090-13	39 39	TY TY	MG MG	O V
0.650770-01 0.	-0.466140-03 0.	0.817650-02 -0.643930-11	-0.733990-02 0.213600-07	0.650670-01 0.289470-10	0.480010-03 -0.823990-14	39 39	TZ TZ	MG MG	O V
0.927680-01 0.108420-16	0.651010-01 0.	-0.710400-08 -0.688340-12	0.316810-08 0.213690-08	0.927680-01 0.289600-11	0.651010-01 -0.867360-15	39 39	VE VE	MG MG	O V
0. 0.162630-16	-0.948680-17 0.	0.827970-02 -0.222040-12	0.100190 01 0.100000 01	0.205190-02 0.135820-02	0.140390-04 -0.260210-15	40 40	VL VL	TD TD	O V
0.292270-03 0.510310-03	0.163700-03 -0.214990-03	-0.322000-04 -0.316540-06	0.115610-04 -0.181620-06	0.292310-03 0.510310-03	0.163690-03 0.214990-03	41 41	RX RX	MG MG	O V
0.481550-03 0.176180-14	-0.114970-05 -0.108420-14	0.602000-04 -0.133230-09	0.141370-04 -0.605180-05	0.481620-03 -0.820280-08	0.125180-05 -0.173470-12	41 41	RY RY	MG MG	O V
0.893830-05 0.294540-01	0.206020-04 0.654520-01	-0.112460-03 0.631830-06	0.247580-03 -0.679670-06	0.961370-05 0.294540-01	0.206020-04 0.654520-01	41 41	RZ RZ	MG MG	O V
0.214000-02 0.128750-15	0.475370-02 0.	0.270820-03 -0.643930-11	0.655010-01 0.654380-01	0.227370-02 0.886970-04	0.475420-02 -0.823990-14	41 41	TX TX	MG MG	O V
0. 0.149080-15	-0.108420-15 0.	-0.714620-07 -0.777160-11	-0.828870-08 -0.185130-07	-0.589780-10 -0.250980-10	-0.121170-09 -0.997470-14	41 41	TY TY	MG MG	O V
0.653680-01 0.813150-16	-0.155620-03 0.	0.827490-02 -0.532910-11	0.192280-02 -0.185220-07	0.653770-01 -0.251090-10	0.189650-03 -0.650520-14	41 41	TZ TZ	MG MG	O V
0.294420-01 0.745390-17	0.654030-01 0.	-0.714660-08 -0.688340-12	-0.829400-09 -0.185130-08	0.294420-01 -0.250970-11	0.654030-01 -0.867360-15	41 41	VE VE	MG MG	O V

END-OF-DATA ENCOUNTERED ON SYSTEM INPUT FILE.